

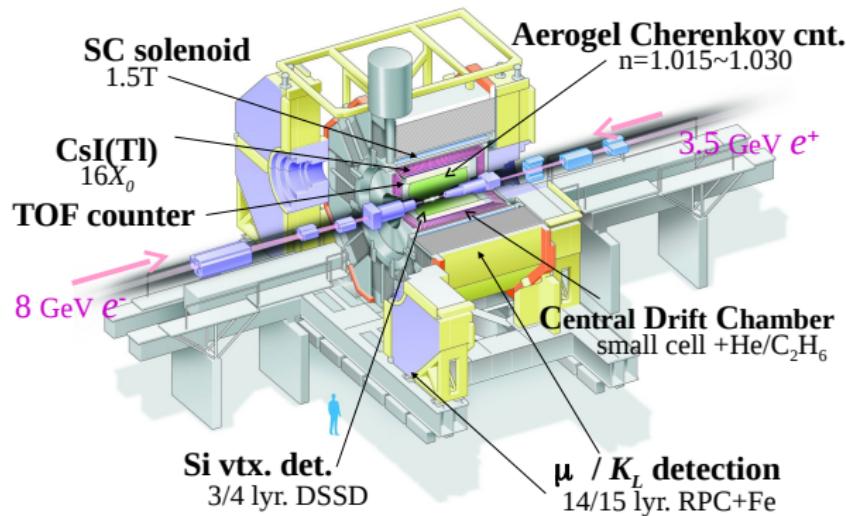
Overview of Belle results

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Belle Detector



The Belle experiment operated at the asymmetric e^+e^- collider KEKB.
Data samples:

- $\Upsilon(4S)$: 711 fb^{-1} , $772 \times 10^6 B\bar{B}$ pairs.
- All energies: 980 fb^{-1} .

Outline

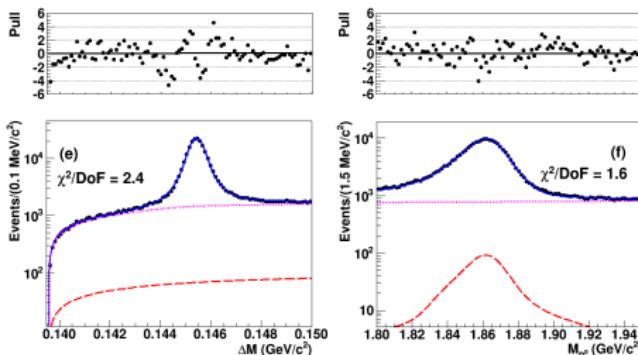
Belle results from 2017 and 2018 are included into this talk. There are following categories by physics subject:

1. CP violation.
2. Standard model tests.
3. Spectroscopy.
4. Transitions between quarkonium states.
5. Initial state radiation processes.
6. τ physics.
7. $\gamma\gamma$ processes.
8. Branching fraction and cross section measurements.

CP violation

Definitions:

- Momenta: $C_T = p_1 \cdot (p_2 \times p_3)$; for D^0 (C_T): 1 = K_S^0 , 2 = π^+ , 3 = π^- ; for \bar{D}^0 (\bar{C}_T): 1 = K_S^0 , 2 = π^- , 3 = π^+ .
- CP asymmetries: $A_T = \frac{\Gamma(C_T > 0) - \Gamma(C_T < 0)}{\Gamma(C_T > 0) + \Gamma(C_T < 0)}$; \bar{A}_T : same with $C_T \rightarrow -\bar{C}_T$.
- T -odd asymmetry: $a_{CP}^{T-\text{odd}} = \frac{1}{2}(A_T - \bar{A}_T)$. This quantity is not affected by final-state interactions.



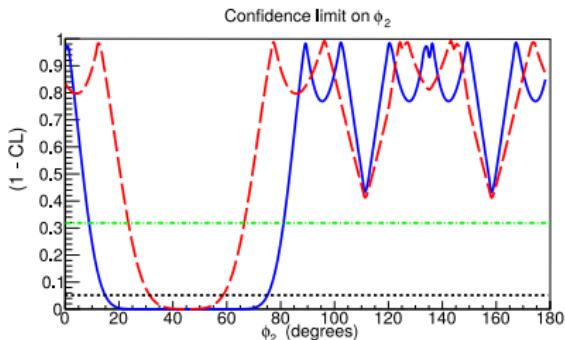
Result:

$$a_{CP}^{T-\text{odd}} = (-0.28 \pm 1.38^{+0.23}_{-0.76}) \times 10^{-3}.$$

Consistent with no CPV.
SM: $\mathcal{O}(10^{-3})$ in D decays.

Data set: $M_{D^*} - M_D$, M_D

- Definition: $A_{CP} = \frac{\Gamma(\bar{B}^0 \rightarrow \pi^0\pi^0) - \Gamma(B^0 \rightarrow \pi^0\pi^0)}{\Gamma(\bar{B}^0 \rightarrow \pi^0\pi^0) + \Gamma(B^0 \rightarrow \pi^0\pi^0)}$.
- Flavor of B is determined by tagging the second B meson by its decay products. The final state is neutral \implies the vertex cannot be determined \implies additional dilution of B^0 and \bar{B}^0 due to mixing.
- The asymmetry is extracted from an unbinned maximum likelihood fit, parameter space: ΔE , M_{bc} , classifier (Fisher discriminant) output, “flavor charge” (1 for B^0 and -1 for \bar{B}^0).
- Angle α/ϕ_2 : from combination with other results ($B^0 \rightarrow \pi^+\pi^-$ and $B^+ \rightarrow \pi^+\pi^0$), method: see PRL **65**, 3381 (1990).



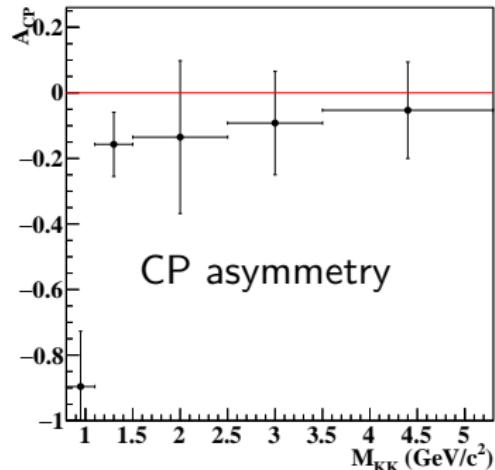
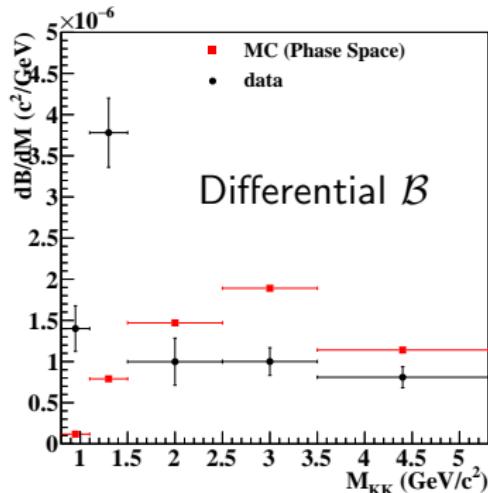
Results:

$$A_{CP} = 0.14 \pm 0.36 \pm 0.10,$$

$$\mathcal{B}(B^0 \rightarrow \pi^0\pi^0) =$$

$$(1.31 \pm 0.19 \pm 0.19) \times 10^{-6},$$

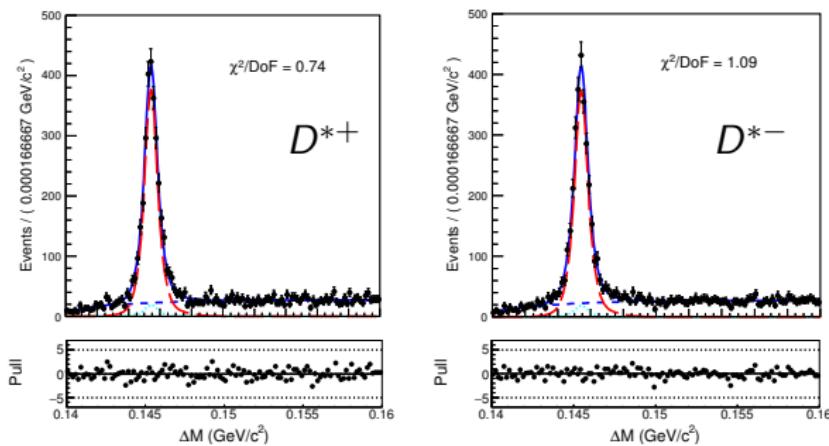
ϕ_2 constraint: see figure (blue solid line; red dashed line is the old constraint from PRD **88**, 092003 (2013)).



Results: $A_{CP} = -0.170 \pm 0.073 \pm 0.017$,
 $\mathcal{B}(B^+ \rightarrow K^+ K^- \pi^+) = (5.38 \pm 0.40 \pm 0.35) \times 10^{-6}$.

- Excess in $K^+ K^-$ low-mass spectrum.
- Large A_{CP} for $M_{K^+ K^-} < 1.1 \text{ GeV}/c^2$ (non-zero: 4.8σ). FSI?

This confirms results by BaBar [PRL 99, 221801 (2007)] and LHCb [PRL 112, 011801 (2014), PRD 90, 112004 (2014)].



Data set: $M_{D^*} - M_D$

Simultaneous fit to $D^{*+} \rightarrow D^0(\rightarrow K_S^0 K_S^0)\pi^+$ and

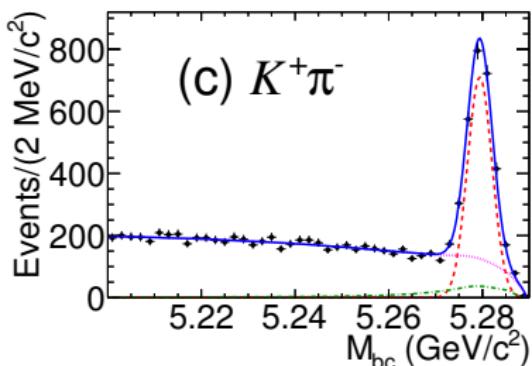
$D^{*-} \rightarrow \bar{D}^0(\rightarrow K_S^0 K_S^0)\pi^-$. Results:

$$\mathcal{B}(D^0 \rightarrow K_S^0 K_S^0) = (1.321 \pm 0.023 \pm 0.036 \pm 0.030) \times 10^{-4},$$

$$A_{CP} = (-0.02 \pm 1.53 \pm 0.02 \pm 0.17) \times 10^{-2}.$$

Third uncertainties are that of $\mathcal{B}(D^0 \rightarrow K_S^0 \pi^0)$.

Theoretical prediction [PRD 92, 054036 (2016)]: $A_{CP} < 1.1 \times 10^{-2}$.



Reconstruction: $K^*(892)^0$: $K_S^0\pi^0$, $K^-\pi^+$; $K^*(892)^+$: $K^+\pi^0$, $K_S^0\pi^+$.
Mass:

$|M_{K\pi} - m_{K^*(892)}| < 75 \text{ MeV}/c^2$,
events with $M_{K\pi} < 2 \text{ GeV}/c^2$ are used to calculate the non- $K^*(892)$ contribution.

Data for $B^+ \rightarrow K^*(892)(\rightarrow K^-\pi^+)\gamma$

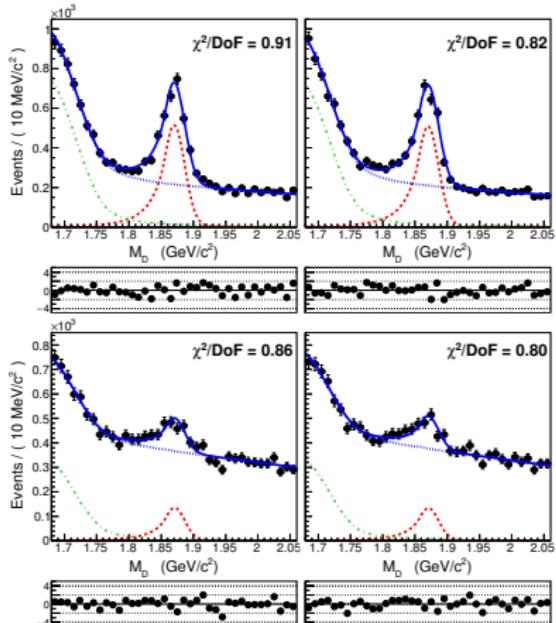
- Isospin asymmetry (3.1σ evidence):

$$\Delta_{0+} = \frac{\Gamma(B^0 \rightarrow K^{*0}\gamma) - \Gamma(B^+ \rightarrow K^*\gamma)}{\Gamma(B^0 \rightarrow K^{*0}\gamma) + \Gamma(B^+ \rightarrow K^*\gamma)} = (6.2 \pm 1.5 \pm 0.6 \pm 1.2) \times 10^{-2}.$$

The fourth uncertainty is that of $\frac{\mathcal{B}(\Upsilon(4S) \rightarrow B^+ B^-)}{\mathcal{B}(\Upsilon(4S) \rightarrow B^0 \bar{B}^0)}$.

- Difference of CP asymmetries (consistent with 0):

$$\Delta A_{CP} = A_{CP}(B^+ \rightarrow K^{*+}\gamma) - A_{CP}(B^0 \rightarrow K^{*0}\gamma) = (2.4 \pm 2.8 \pm 0.5) \times 10^{-2}$$



The D mass for
 $p_{D^*}^* > 2.95 \text{ GeV}/c$ (top),
 $2.50 < p_{D^*}^* < 2.95 \text{ GeV}/c$
(bottom), D^+ (left), D^- (right).

Raw asymmetry:

$A_{\text{raw}}^{\pi\pi} = A_{CP}^{\pi\pi} + A_{FB} + A_\epsilon^{\pi^\pm}$, A_{FB} - forward-backward, $A_\epsilon^{\pi^\pm}$ - pion detection asymmetry; a normalization channel $D^+ \rightarrow K_S^0\pi^+$ is used to correct for them.

Tagged (D^+ from $D^{*+} \rightarrow D^+\pi^0$) and untagged samples (no D^{*+} candidates are found) are used. The samples are fitted separately. Results:

$$A_{\text{raw}}^{\pi\pi} - A_{\text{raw}}^{K\pi} = (0.81 \pm 1.97 \pm 0.19) \times 10^{-2}$$

(tagged), $(4.02 \pm 1.61 \pm 0.32) \times 10^{-2}$
(untagged).

After that, the results are combined, and $A_{CP}(D^+ \rightarrow K_S^0\pi^+)$ is added. The final result is:

$$A_{CP} = (2.31 \pm 1.24 \pm 0.23) \times 10^{-2}.$$

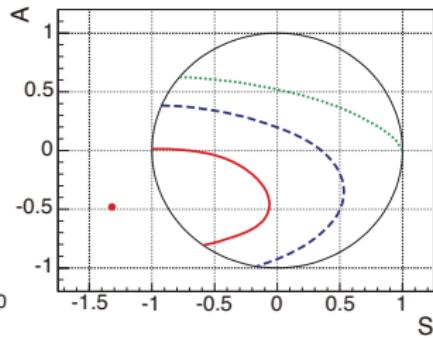
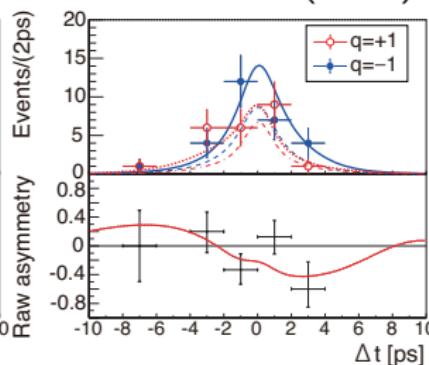
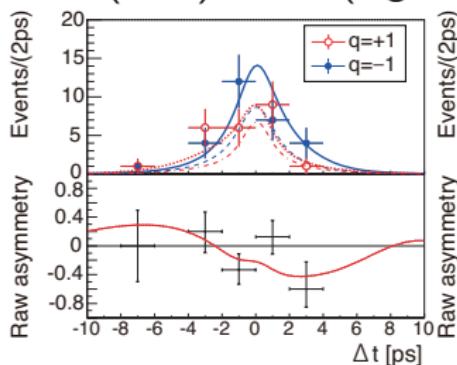
$$\mathcal{P}(\Delta t) = \frac{e^{-|\Delta t|/\tau_{B^0}}}{4\tau_{B^0}} \{1 + q [\mathcal{S} \sin(\Delta m_d \Delta t) + \mathcal{A} \cos(\Delta m_d \Delta t)]\}, \quad (1)$$

\mathcal{S} - mixing-induced, \mathcal{B} - direct CP violation parameters, q - flavor charge.
 Fit to $\mathcal{P}(t)$ modified to incorporate incorrect flavor assignment. Results:
 $\mathcal{S} = -1.32 \pm 0.77 \pm 0.36$, $\mathcal{A} = -0.48 \pm 0.41 \pm 0.07$.

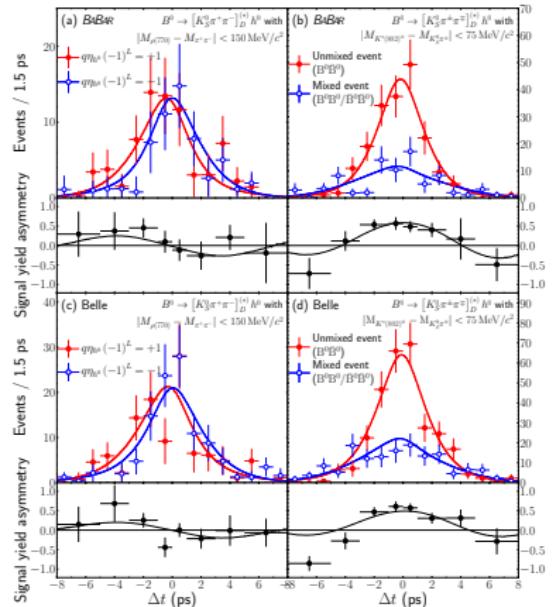
The central point is out of the physical region $\mathcal{S}^2 + \mathcal{A}^2 \leq 1$.

Theoretical prediction (SM):

$\mathcal{S} \sim -(2-3) \times 10^{-2}$ (e.g. PRD **73**, 114022 (2006), PLB **642**, 478 (2006)).



This analysis has been performed by Belle and BaBar.



$$D^{*0} \rightarrow D^0 \pi^0, D^0 \rightarrow K_S^0 \pi^+ \pi^-, h^0 = \pi^0, \eta, \omega.$$

Decay modes (sufficient signal yield):

$$B^0 \rightarrow D^0 \pi^0, D^0 \eta, D^0 \omega, D^{*0} \pi^0, D^0 \eta.$$

The amplitude of $D^0 \rightarrow K_S^0 \pi^+ \pi^-$ is determined by performing a Dalitz analysis of this decay for continuum D mesons. After that, a time-dependent Dalitz analysis is performed (D decay amplitude is fixed, only $\sin 2\beta$ and $\cos 2\beta$ are free). Results:
 $\sin 2\beta = 0.80 \pm 0.14 \pm 0.06 \pm 0.03$,
 $\cos 2\beta = 0.91 \pm 0.22 \pm 0.09 \pm 0.07$.

The fourth uncertainties are those of the $D^0 \rightarrow K_S^0 \pi^+ \pi^-$ amplitude model. Direct angle measurement (β is a fit parameter):
 $\beta = (22.5 \pm 4.4 \pm 1.2 \pm 0.6)$ degrees.

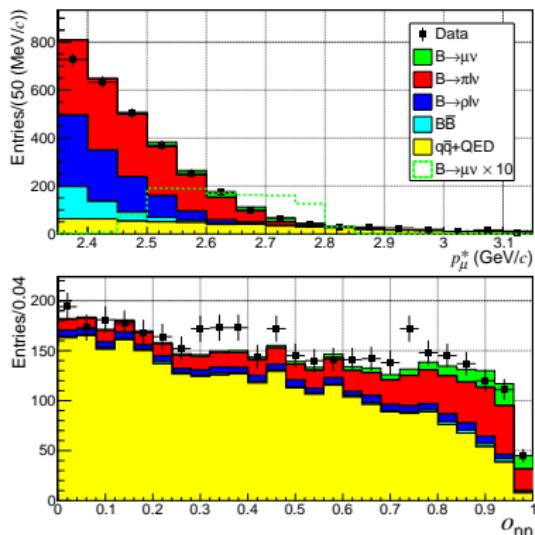
The second minimum ($\beta' = \pi/2 - \beta$) is excluded at 7.3σ level.

Standard model tests

SM expectation:

$$\mathcal{B}(B^- \rightarrow \ell^- \bar{\nu}_\ell) = \frac{G_F^2 m_B m_\ell^2}{8\pi} \left(1 - \frac{m_\ell^2}{m_B^2}\right)^2 f_B^2 |V_{ub}|^2 \tau_B \quad (2)$$

$f_B = 0.186 \pm 0.004$ (from lattice) is the B decay constant.



The result is

$$\mathcal{B}(B^- \rightarrow \mu^- \bar{\nu}_\mu) = (3.80 \pm 0.31) \times 10^{-7}$$

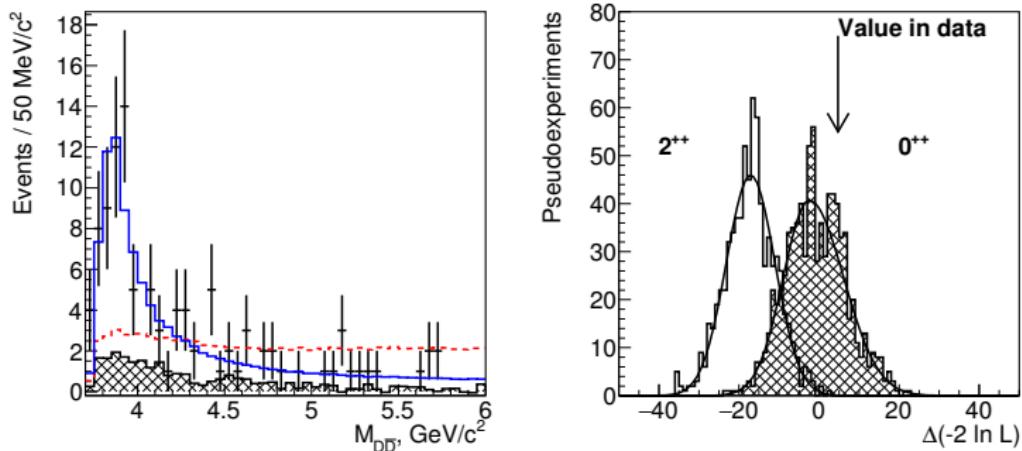
Initial selection uses accompanying B meson candidate ($M_{bc} > 5.1 \text{ GeV}/c^2$, $-3 < \Delta E < 2 \text{ GeV}$).

Number of signal events is determined by fitting the distribution of neural network output and muon momentum. Result:

$$\mathcal{B}(B^- \rightarrow \mu^- \bar{\nu}_\mu) = (6.46 \pm 2.22 \pm 1.60) \times 10^{-7}$$

Significance: 2.4σ .

Spectroscopy



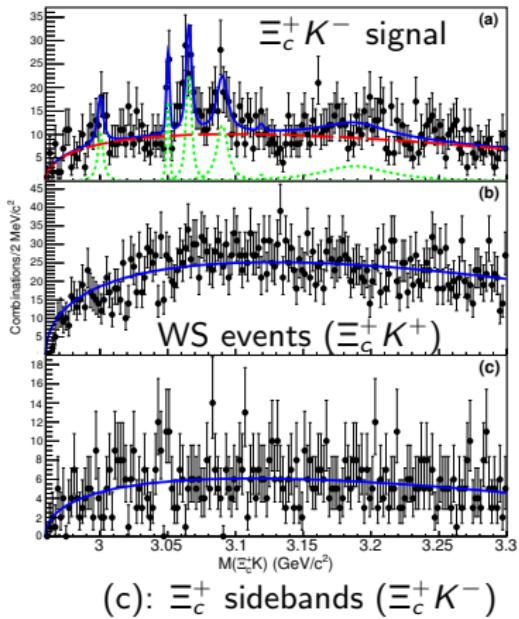
Full 6-dimensional amplitude analysis.

A new state $X^*(3860)$ is observed, 6.5σ including systematic error.

Parameters: $M = 3862^{+26+40}_{-32-13}$ MeV/ c^2 , $\Gamma = 201^{+154+88}_{-67-82}$ MeV.

$J^{PC} = 0^{++}$ is preferred over 2^{++} at 2.5σ .

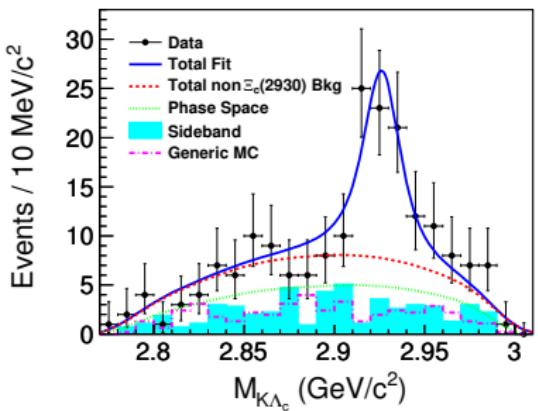
Mass, mass difference with the $\chi_{c2}(2P)$, J^{PC} , decay mode, S -wave production \implies the $X^*(3860)$ is a better $\chi_{c0}(2P)$ candidate than the $X(3915)$ that is observed in $B \rightarrow J/\psi \omega K$ and $\gamma\gamma \rightarrow J/\psi \omega$.



- LHCb observed 5 new narrow Ω_c states in PRL **118**, 182001 (2017).
- Belle searched for continuum production of these states in the same $\Xi_c^+ K^-$ decay mode.
- 4 of 5 states are observed, 2 of them with significance of $> 5\sigma$.
- The $\Omega_c(3119)$ is not observed as well as the $\Omega_c(3188)$ (possible wide state at higher mass from the same LHCb analysis).

Confirmed states:

State	$\Omega_c(3000)$	$\Omega_c(3050)$	$\Omega_c(3066)$	$\Omega_c(3090)$
Yield	37.7 ± 11.0	28.2 ± 7.7	81.7 ± 13.9	86.6 ± 17.4
Significance	3.9σ	4.6σ	7.2σ	5.7σ
LHCb Mass	$3000.4 \pm 0.2 \pm 0.1$	$3050.2 \pm 0.1 \pm 0.1$	$3065.5 \pm 0.1 \pm 0.3$	$3090.2 \pm 0.3 \pm 0.5$
Belle Mass	$3000.7 \pm 1.0 \pm 0.2$	$3050.2 \pm 0.4 \pm 0.2$	$3064.9 \pm 0.6 \pm 0.2$	$3089.3 \pm 1.2 \pm 0.2$



- The $\Xi_c(2930)^0$ has previously been reported by BaBar in PRD **77**, 031101 (2008) in the same process, but the significance has not been specified \implies its status is unclear.
- Unbinned simultaneous extended maximum likelihood fit to the signal and Λ_c sidebands. The fit is 1-dimensional ($\Lambda_c^+ K^-$).

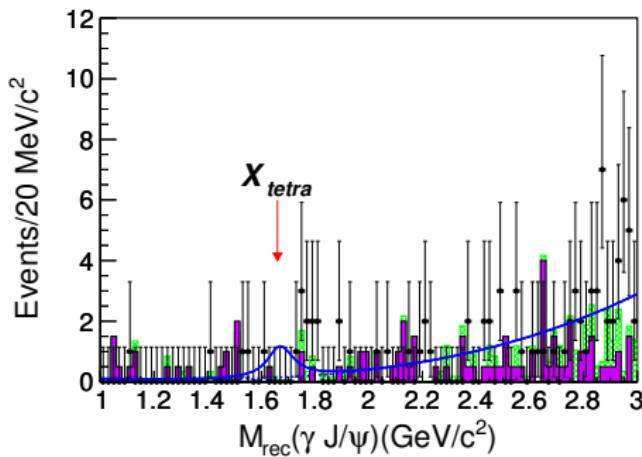
The $\Xi_c(2930)^0$ is observed with 5.1σ global significance.

Parameters: $M = 2928.9 \pm 3.0^{+0.9}_{-12.0} \text{ MeV}/c^2$, $\Gamma = (19.5 \pm 8.4^{+5.9}_{-7.9}) \text{ MeV}$.

$$\mathcal{B}(B^- \rightarrow \Lambda_c^+ \bar{\Lambda}_c^- K^-) = (4.80 \pm 0.43 \pm 0.60) \times 10^{-4}$$

$$\mathcal{B}(B^- \rightarrow \Xi_c(2930)^0 \bar{\Lambda}_c^-) \times \mathcal{B}(\Xi_c(2930)^0 \rightarrow \Lambda_c^+ K^-) = (1.73 \pm 0.45 \pm 0.21) \times 10^{-4}$$

$$\mathcal{B}(B^- \rightarrow Y(4660) K^-) \times \mathcal{B}(Y(4660) \rightarrow \Lambda_c^+ \bar{\Lambda}_c^-) < 1.2 \times 10^{-4} \text{ (90% C.L.)}$$

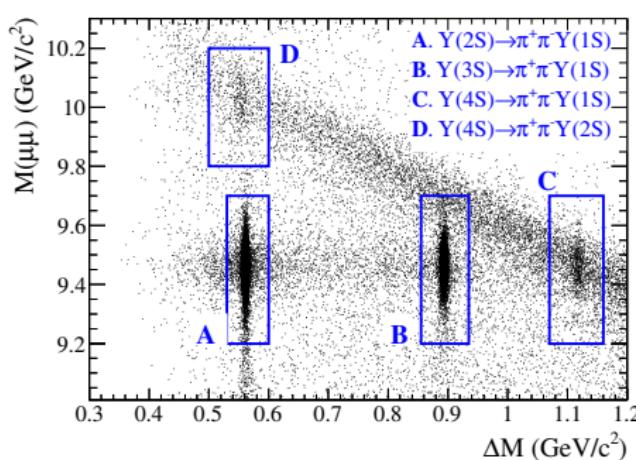


Search for light tetraquarks in $\Upsilon(1S, 2S) \rightarrow \chi_{c1} + X$,
 $\Upsilon(1S, 2S) \rightarrow f_1(1285) + X$,
 $\chi_{b1} \rightarrow J/\psi + X$, $\chi_{b1} \rightarrow \omega X$ (the choice of the channels is such that the X state can have $J^{PC} = 0^{--}$) by fits of the recoil mass spectra.

No signal is found, upper limits are calculated as functions of the X mass.
Inclusive branching fractions:

$$\begin{aligned}\mathcal{B}(\Upsilon(1S) \rightarrow f_1(1285) + \dots) &= (4.6 \pm 2.8 \pm 1.3) \times 10^{-3}, \\ \mathcal{B}(\Upsilon(2S) \rightarrow f_1(1285) + \dots) &= (2.2 \pm 1.5 \pm 0.6) \times 10^{-3}, \\ \mathcal{B}(\chi_{b2} \rightarrow J/\psi + \dots) &= (1.50 \pm 0.34 \pm 0.22) \times 10^{-3}, \\ \mathcal{B}(\chi_{b1} \rightarrow \omega + \dots) &= (4.9 \pm 1.3 \pm 0.6) \times 10^{-2}, \\ \mathcal{B}(\chi_{b0}[\chi_{b1}] \rightarrow J/\psi + \dots) &< 2.3[1.1] \times 10^{-3} \text{ (90% C.L.)}.\end{aligned}$$

Transitions between quarkonium states



$$\Upsilon(4S) \rightarrow \Upsilon(1S, 2S) \pi^+ \pi^- ,$$

$$\Upsilon(4S) \rightarrow \Upsilon(1S) \eta \text{ with}$$

$$\Upsilon(1S, 2S) \rightarrow \mu^+ \mu^- .$$

$$\Delta M = M_{\mu^+ \mu^- \pi^+ \pi^-} - M_{\mu^+ \mu^-}$$

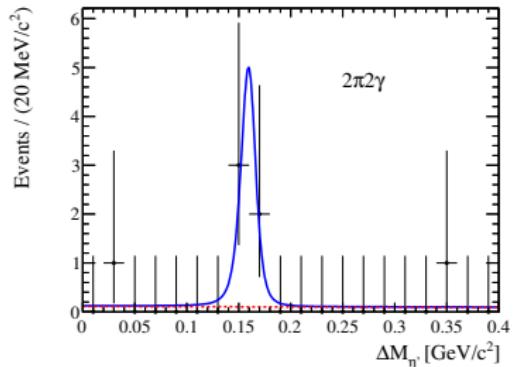
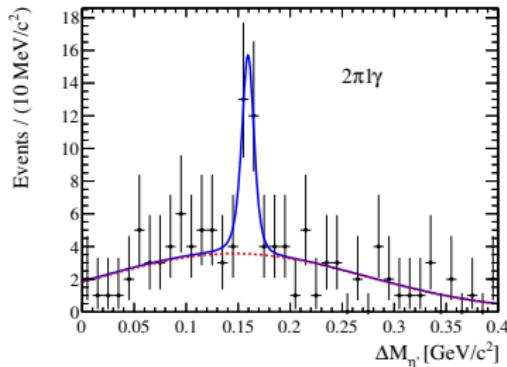
$$\mathcal{B}(\Upsilon(4S) \rightarrow \Upsilon(1S) \pi^+ \pi^-) = (8.2 \pm 0.5 \pm 0.4) \times 10^{-5},$$

$$\mathcal{B}(\Upsilon(4S) \rightarrow \Upsilon(2S) \pi^+ \pi^-) = (7.9 \pm 1.0 \pm 0.4) \times 10^{-5},$$

$$\mathcal{B}(\Upsilon(4S) \rightarrow \Upsilon(1S) \eta) = (1.70 \pm 0.23 \pm 0.08) \times 10^{-4}.$$

Transition from $\Upsilon(4S)$ via η is not suppressed despite of the b quark spin flip. Confirms similar BaBar result (PRD **78**, 112002).

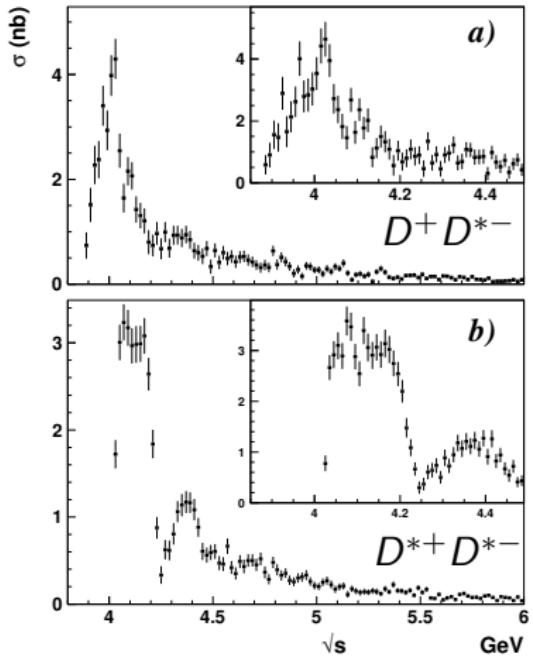
ISR cross sections at $\Upsilon(4S)$: $\sigma_{\text{ISR}}(\Upsilon(2S)) = (17.36 \pm 0.19 \pm 0.69) \text{ pb}$,
 $\sigma_{\text{ISR}}(\Upsilon(3S)) = (28.9 \pm 0.5 \pm 1.3) \text{ pb}$.



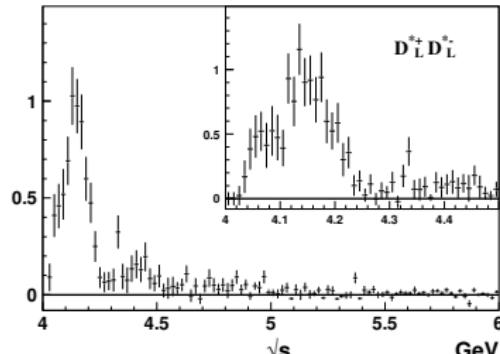
- Reconstruction: $\eta' \rightarrow \rho^0\gamma$, $\pi^+\pi^-\eta(\rightarrow\gamma\gamma)$; $\Upsilon(1S) \rightarrow \mu^+\mu^-$.
- η : full reconstruction (left figure) or 1 photon only (right figure).
- Background PDF: linear (full reconstruction), Gaussian (partial reconstruction).
- Significance (combined by simultaneous fit): 5.7σ (with systematic uncertainty).
- Branching: $\mathcal{B}(\Upsilon(4S) \rightarrow \eta'\Upsilon(1S)) = (3.43 \pm 0.88 \pm 0.21) \times 10^{-5}$.

Initial state radiation processes

Cross sections:

Partial reconstruction: D^0 from D^{*+} is not reconstructed.Distributions of the D^* helicity angles are measured (for each mass bin):

- $D^{*+}D^{*-}$: no parameters.
- $D^{*+}D^{*-}$: distribution depends on 3 cross sections σ_{LL} , σ_{TL} , σ_{TT} (L : longitudinal, $\lambda = 0$; T : transverse, $\lambda = \pm 1$). Example result:

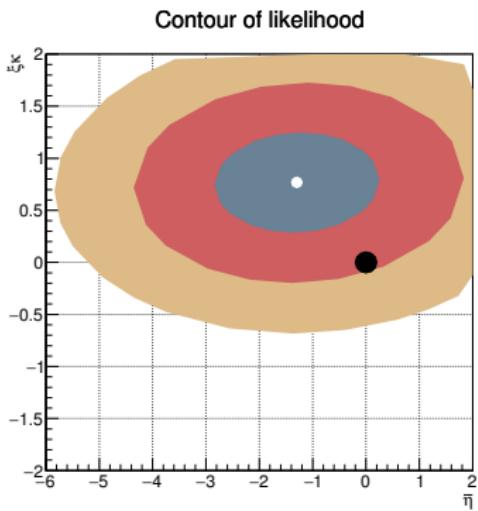


τ physics

$$\bar{\eta} = \left|g_{RL}^V\right|^2 + \left|g_{LR}^V\right|^2 + \frac{1}{8} \left(\left|g_{RL}^S + 2g_{RL}^T\right|^2 + \left|g_{LR}^S + 2g_{LR}^T\right|^2 \right) + 2 \left(\left|g_{RL}^T\right|^2 + \left|g_{LR}^T\right|^2 \right),$$

$$\xi\kappa = \left|g_{RL}^V\right|^2 - \left|g_{LR}^V\right|^2 + \frac{1}{8} \left(\left|g_{RL}^S + 2g_{RL}^T\right|^2 - \left|g_{LR}^S + 2g_{LR}^T\right|^2 \right) + 2 \left(\left|g_{RL}^T\right|^2 - \left|g_{LR}^T\right|^2 \right).$$

In SM: $g_{LL}^V = 1$, other couplings = 0 $\implies \bar{\eta} = \xi\kappa = 0$.



The parameters $\bar{\eta}$ and $\xi\kappa$ are extracted by simultaneous fit to differential cross sections of $\tau^- \rightarrow \ell^- \nu_\tau \bar{\nu}_\ell \gamma$. Results:

$$\bar{\eta} = -1.3 \pm 1.5 \pm 0.8,$$

$$\xi\kappa = 0.5 \pm 0.4 \pm 0.2.$$

Branching fractions with $E_\gamma > 10$ MeV:

$$\mathcal{B}(\tau^- \rightarrow e^- \nu_\tau \bar{\nu}_e \gamma) = (1.79 \pm 0.02 \pm 0.10) \times 10^{-2},$$

$$\mathcal{B}(\tau^- \rightarrow \mu^- \nu_\tau \bar{\nu}_\mu \gamma) = (3.63 \pm 0.02 \pm 0.15) \times 10^{-3}.$$

Ratio: $4.95 \pm 0.06 \pm 0.20$ (prediction: 4.605).

Measured parameters:

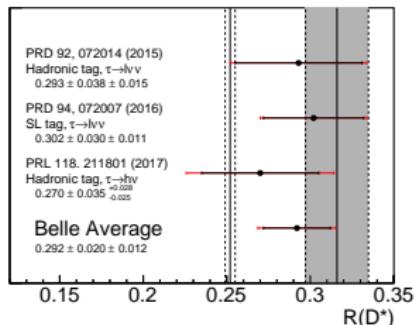
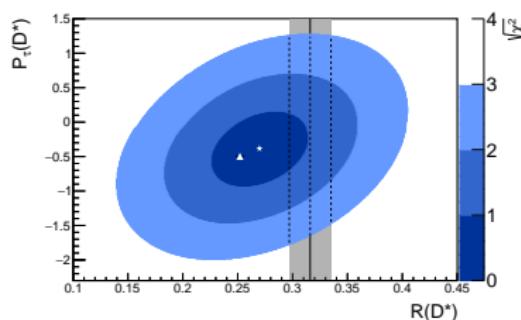
- $R(D^*) = \frac{\mathcal{B}(\bar{B} \rightarrow D^* \tau^- \bar{\nu}_\tau)}{\mathcal{B}(\bar{B} \rightarrow D^* \ell^- \bar{\nu}_\ell)}$.

For Belle and BaBar, $\ell = e + \mu$ average; for LHCb, $\ell = \mu$.

- $P_\tau(D^*) = \frac{\Gamma^+ - \Gamma^-}{\Gamma^+ + \Gamma^-}$,

where Γ^+ , Γ^- - decay width with $\lambda_\tau = +1/2$ and $-1/2$.

Results: $R(D^*) = 0.270 \pm 0.035^{+0.028}_{-0.025}$, $P_\tau(D^*) = -0.38 \pm 0.51^{+0.21}_{-0.16}$.



Average of 3 Belle measurements: $R(D^*) = 0.292 \pm 0.020 \pm 0.012$.

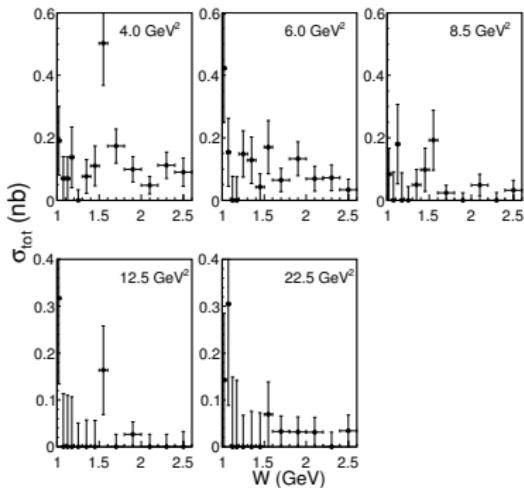
It is 1.7σ larger than SM prediction (0.252 ± 0.003).

New world average is 3.5σ larger.

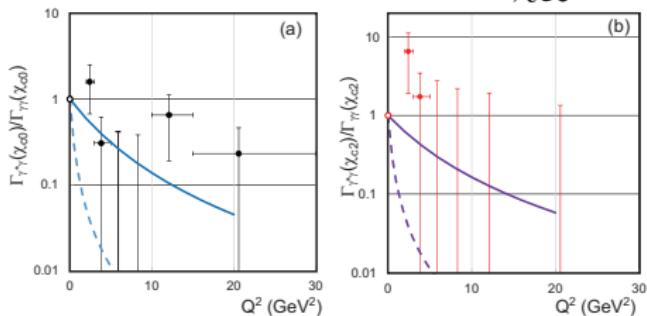
$\gamma\gamma$ processes

Process: $e^+e^- \rightarrow e^+e^- K_S^0 K_S^0$, one of the final-state e^\pm is detected, the other one is not. The cross sections depend on $Q^2 = -m_{\gamma^*}^2$.

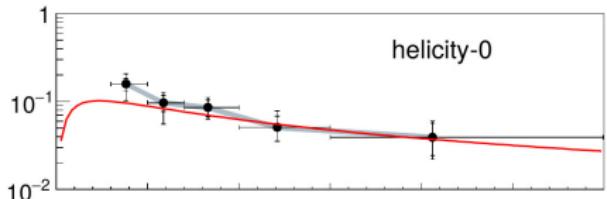
Total cross sections in bins of Q^2 (bin center is specified on the histograms).



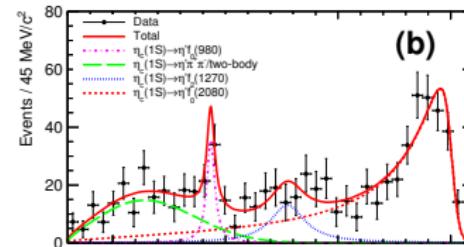
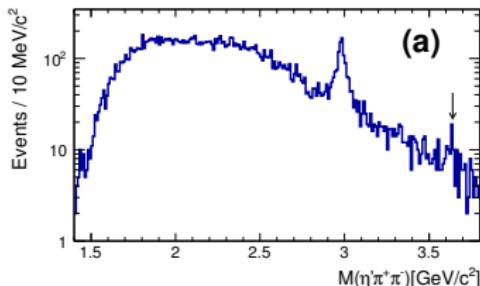
Differential cross sections for χ_{c0} and χ_{c2} .



Also, transition form factor of the $f'_2(1525)$ is measured for all helicity components (0, 1, 2).



Reconstruction: no e^\pm tagging, selection by low p_t .

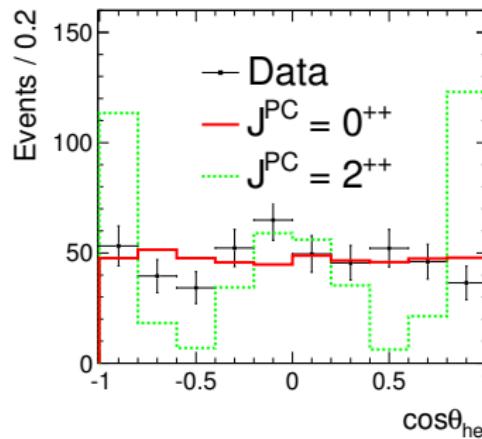


Observed:

- $\eta_c(2S) \rightarrow \eta'\pi^+\pi^-$ (5.5σ).
- $\eta_c(1S) \rightarrow \eta'f_0(2080)$ (20σ).

Measured:

- $f_0(2080)$: $J^{PC} = 0^{++}$
(exclusion of 2^{++} : 11σ).
- Parameters of η_c , $\eta_c(2S)$, $f_0(2080)$.



Branching fraction and cross section measurements

Branching fractions:

- $B^+ \rightarrow \eta\ell^+\nu_\ell$, $B^+ \rightarrow \eta'\ell^+\nu_\ell$: PRD **96**, 091102 (2017)
- $\Lambda_c^+ \rightarrow p\phi\pi^0$, $\Lambda_c^+ \rightarrow pK^-\pi^+\pi^0$: PRD **96**, 051102 (2017)
- $B^+ \rightarrow X_{c\bar{c}}K^+$, $B^+ \rightarrow \bar{D}^{(*)0}\pi^+$: PRD **97**, 012005 (2018)
- Ω_c hadronic decays: PRD **97**, 032001 (2018)
- $\Lambda_c^+ \rightarrow \Sigma^+\pi^+\pi^-$, $\Lambda_c^+ \rightarrow \Sigma^+\pi^0\pi^0$, $\Lambda_c^+ \rightarrow \Sigma^0\pi^+\pi^0$: arXiv:1802.03421
- $B \rightarrow D^{(*)}\pi\ell\nu$: arXiv:1803.06444

Cross sections:

- $e^+e^- \rightarrow$ hyperons, charmed baryons: PRD **97**, 072005 (2018)
- $e^+e^- \rightarrow hhX$ ($h = \pi, K$): PRD **96**, 032005 (2017)
- $e^+e^- \rightarrow \eta\Upsilon$, $e^+e^- \rightarrow \eta h_b$: arXiv:1803.03225

Thank you for attention!