

Holographic anisotropic background with confinement-deconfinement phase transition

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based on I.Ya. Aref'eva and K. Rannu,
JHEP (2018), arXiv:1802.05652 [hep-th]

Motivation

Holography needs:

- **Natural** theory with metric as solution of EOM.
- Non-zero chemical potential $\mu \neq 0$.

I.Ya. Aref'eva “**Holography for HIC at LHC and NICA**”
arXiv:1612.08928

$$ds^2 = L^2 \frac{b(z)}{z^2} \left[-g(z)dt^2 + dx^2 + z^{2-\frac{2}{\nu}} (dy_1^2 + dy_2^2) + \frac{dz^2}{g(z)} \right]$$

I.A., K.R. arXiv:1802.05652

AGG solution

$$ds^2 = \frac{L^2}{z^2} \left[-g(z)dt^2 + dx^2 + z^{2-\frac{2}{\nu}} (dy_1^2 + dy_2^2) + \frac{dz^2}{g(z)} \right]$$

$\mu = 0, b(z) = 1, g(z) = 1$ Aref'eva, Golubtsova JHEP **1504** 011

$\mu = 0, b(z) = 1, g(z) \neq 1$ Aref'eva, Golubtsova, Gourgoulhon JHEP **1609** 142

Since phenomenology requires $\mu \neq 0, b \neq 1, \nu \neq 1$ Aref'eva 1612.08928

OUR GOAL



Generalizations:



$\mu \neq 0$



$b(z) \neq 1$

Action in Einstein frame, metric ansatz

$$S = \frac{1}{16\pi G_5} \int d^5x \sqrt{-g} \times$$
$$\times \left[R - \frac{f_1(\phi)}{4} F_{(1)}^2 - \frac{f_2(\phi)}{4} F_{(2)}^2 - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) \right]$$

$$A_\mu^{(1)} = A_t(z) \delta_\mu^0 \quad F^{(2)} = q dy^1 \wedge dy^2 \quad \phi = \phi(z)$$

$$ds^2 = L^2 \frac{b(z)}{z^2} \left[-g(z) dt^2 + dx^2 + z^{2-\frac{2}{\nu}} (dy_1^2 + dy_2^2) + \frac{dz^2}{g(z)} \right]$$

Boundary conditions:

$$g(0) = 1, \quad g(z_h) = 0 \quad A_t(0) = \mu, \quad A_t(z_h) = 0 \quad \phi(z_h) = 0$$

$\nu = 1$: M.-W. Li, Y. Yang arXiv:1703.09184

Equations of motion

$$\phi'' + \phi' \left(\frac{g'}{g} + \frac{3b'}{2b} - \frac{\nu + 2}{\nu z} \right) + \frac{z^2 A_t'^2}{2bg} \frac{\partial f_1}{\partial \phi} - \frac{q^2 z^{-2 + \frac{4}{\nu}}}{2bg} \frac{\partial f_2}{\partial \phi} - \frac{b}{z^2 g} \frac{\partial V}{\partial \phi} = 0$$

$$A_t'' + A_t' \left(\frac{b'}{2b} + \frac{f_1'}{f_1} - \frac{2 - \nu}{\nu z} \right) = 0$$

$$\text{1-st: } g'' + g' \left(\frac{3b'}{2b} - \frac{1}{z} - \frac{2}{\nu z} \right) - \frac{z^2}{b} f_1 A_t'^2 = 0$$

$$\text{2-nd: } b'' - \frac{3b'^2}{2b} + \frac{2b'}{z} - \frac{4b}{3\nu z^2} \left(1 - \frac{1}{\nu} \right) + \frac{b}{3} \phi'^2 = 0$$

$$\text{3-rd: } 2g' \left(1 - \frac{1}{\nu} \right) + g \left(1 - \frac{1}{\nu} \right) \left(\frac{3b'}{b} - \frac{4}{z} - \frac{4}{\nu z} \right) + \frac{q^2 z^{-1 + \frac{4}{\nu}}}{b} f_2 = 0$$

$$\begin{aligned} \text{4-th: } & -V - \frac{z^4}{2b^2} A_t'^2 f_1 - \frac{3z^2 b' g'}{2b^2} - \frac{3z^2 g b'^2}{b^3} + \frac{9z g b'}{2\nu b^2} + \frac{15z g b'}{2b^2} + \\ & + \frac{z g'}{\nu b} + \frac{2z g'}{b} + \frac{z^2 g \phi'^2}{2b} - \frac{8g}{\nu b} - \frac{4g}{b} = 0 \end{aligned}$$

General anisotropic solution

$$b(z) = \exp P(z), \nu \neq 1$$

$$f_1 = z^{-2 + \frac{2}{\nu}}$$

$$A_t = \tilde{\mu} \int_z^{z_h} e^{-\frac{P(\xi)}{2}} \xi d\xi, \quad \tilde{\mu} = \frac{\mu}{\int_0^{z_h} e^{-\frac{P(\xi)}{2}} \xi d\xi}$$

$$g = 1 + \tilde{\mu}^2 \int_0^z e^{-\frac{3P(\xi)}{2}} \left(\int_0^\xi e^{-\frac{P(\chi)}{2}} \chi d\chi \right) \xi^{1 + \frac{2}{\nu}} d\xi -$$
$$- \frac{1 + \tilde{\mu}^2 \int_0^{z_h} e^{-\frac{3P(\xi)}{2}} \left(\int_0^\xi e^{-\frac{P(\chi)}{2}} \chi d\chi \right) \xi^{1 + \frac{2}{\nu}} d\xi}{\int_0^{z_h} e^{-\frac{3P(\xi)}{2}} \xi^{1 + \frac{2}{\nu}} d\xi} \int_0^z e^{-\frac{3P(\xi)}{2}} \xi^{1 + \frac{2}{\nu}} d\xi$$

$$\phi(z) = C_5 + \int_0^z \sqrt{-3P''(\xi) + \frac{3}{2} P'^2(\xi) - \frac{6}{\xi} P'(\xi) + 4 \frac{\nu - 1}{\xi^2 \nu^2}} d\xi$$

Solution for anisotropic metric ansatz

$$b(z) = \exp(cz^2/2), \nu \neq 1$$

$$f_1 = z^{-2+\frac{2}{\nu}} \quad A_t(z) = \mu \frac{e^{-\frac{cz^2}{4}} - e^{-\frac{cz_h^2}{4}}}{1 - e^{-\frac{cz_h^2}{4}}}$$
$$g(z) = 1 - \frac{z^{2+\frac{2}{\nu}} \mathfrak{G}\left(\frac{3}{4}cz^2\right)}{z_h^{2+\frac{2}{\nu}} \mathfrak{G}\left(\frac{3}{4}cz_h^2\right)} - \frac{\mu^2 cz^{2+\frac{2}{\nu}} e^{\frac{cz_h^2}{2}} \mathfrak{G}(cz^2)}{4 \left(1 - e^{-\frac{cz_h^2}{4}}\right)^2} +$$
$$+ \frac{\mu^2 cz^{2+\frac{2}{\nu}} e^{\frac{cz_h^2}{2}} \mathfrak{G}\left(\frac{3}{4}cz^2\right) \mathfrak{G}(cz_h^2)}{4 \left(1 - e^{-\frac{cz_h^2}{4}}\right)^2 \mathfrak{G}\left(\frac{3}{4}cz_h^2\right)},$$
$$\mathfrak{G}(x) = x^{-1-\frac{1}{\nu}} \gamma\left(1 + \frac{1}{\nu}, x\right) = \sum_{n=0}^{\infty} \frac{(-1)^n x^n}{n!(1+n+\frac{1}{\nu})}$$

$b(z)$ similar to phenom. model of O. Andreev, V.I. Zakharov arXiv:0604204

Scalar field $\phi(z)$ for anisotropic metric

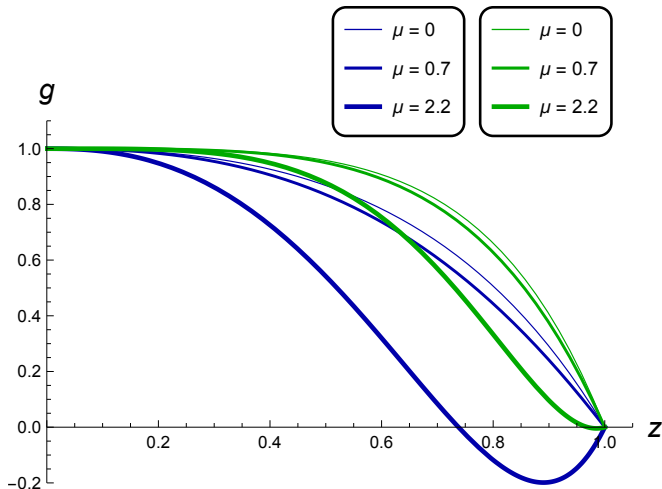
$$b(z) = \exp(cz^2/2), \nu \neq 1$$

$$\begin{aligned} \phi = & \frac{1}{2\sqrt{2}\nu} \left\{ \sqrt{3c^2\nu^2 z^4 - 18c\nu^2 z^2 + 8(\nu - 1)} - \right. \\ & - \sqrt{3c^2\nu^2 z_h^4 - 18c\nu^2 z_h^2 + 8(\nu - 1)} + 2\sqrt{2(\nu - 1)} \ln \left(\frac{z^2}{z_h^2} \right) \\ & - 3\sqrt{3}\nu \ln \left(\frac{\sqrt{3c^2\nu^2 z^4 - 18c\nu^2 z^2 + 8(\nu - 1)} - \sqrt{3}\nu(3 - cz^2)}{\sqrt{3c^2\nu^2 z_h^4 - 18c\nu^2 z_h^2 + 8(\nu - 1)} - \sqrt{3}\nu(3 - cz_h^2)} \right) - \\ & - 2\sqrt{2(\nu - 1)} \times \\ & \left. \times \ln \left(\frac{9c\nu^2 z^2 - 8(\nu - 1) - \sqrt{2(\nu - 1)} \sqrt{3c^2\nu^2 z^4 - 18c\nu^2 z^2 + 8(\nu - 1)}}{9c\nu^2 z_h^2 - 8(\nu - 1) - \sqrt{2(\nu - 1)} \sqrt{3c^2\nu^2 z_h^4 - 18c\nu^2 z_h^2 + 8(\nu - 1)}} \right) \right\} \end{aligned}$$

Stable solution $\Rightarrow c < 0$

Blackening function $g(z)$ for various μ

$z_h = 1, c = -1, \nu = 1, 4.5$



Scalar potential $V(\phi)$ for anisotropic metric

$$b(z) = \exp(cz^2/2), \nu \neq 1$$

$$V(\phi, \mu, \nu) \approx V_0(\nu) - C_7(\mu, \nu)e^{K_1(\nu)\phi} + C_8(\mu, \nu)e^{K_2(\nu)\phi}$$

$$V_0(4.5) = -0.5778$$

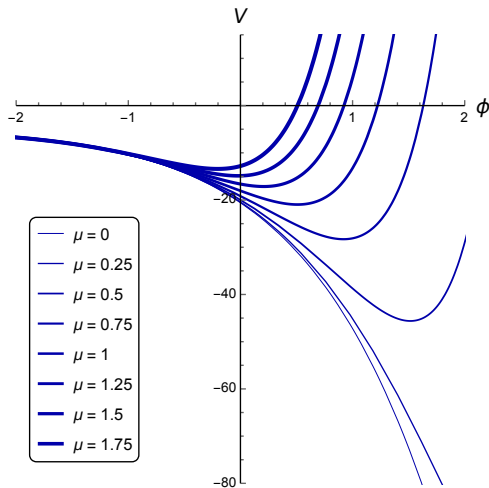
$$K_1(4.5) = 0.7897, \quad K_2(4.5) = 2.0995$$

$$C_7(\mu, 4.5) = 23.0779 + 2.4236\mu^2$$

$$C_8(\mu, 4.5) = 0.0575 + 4.9919\mu^2$$

Scalar potential $V(\phi)$ for various μ

$z_h = 1, c = -1, \nu = 4.5$



Anisotropic solution

$$b(z) = 1, \nu \neq 1, A_t(z) = 0$$

$$g(z) = 1 - \left(\frac{z}{z_h} \right)^{2 + \frac{2}{\nu}}$$

$$f_2(z) = \frac{4z^{-4/\nu}}{q^2} \frac{(\nu - 1)(1 + 3\nu + 2\nu^2)}{\nu^2 (1 + 2\nu)}$$

$$\phi(z) = C_5 \pm 2 \frac{\sqrt{\nu - 1}}{\nu} \log(z)$$

$$V(z) = -2 \frac{(1 + \nu)(1 + 2\nu)}{\nu^2}$$

Corresponds to AGG solution

(I.Ya. Aref'eva, A.A. Golubtsova and E. Gourgoulhon arXiv:1601.06046)

RG flow

$E = E_0 L B$ — energy scale
 $\lambda = e^\phi$ — string coupling function

$z \rightarrow w$:

$$ds^2 = - B(w)^2 g(w) dt^2 + B(w)^2 dx^2 + R(w) (dy_1^2 + dy_2^2) + \frac{dw^2}{g(w)}$$

$$E_0 L = 1$$

\Downarrow

$$X(\phi) = \frac{B}{B'} \frac{\phi'}{3}, \quad Y(\phi) = \frac{1}{4} \frac{g'}{g} \frac{B}{B'}, \quad Z(\phi) = \frac{1}{4} \frac{R'}{R} \frac{B}{B'}$$

$$H_1(\phi) = \frac{A'_t}{B}, \quad H_2(\phi) = \frac{q}{R}$$

*Generalization of: U. Gursoy, E. Kiritsis, L. Mazzanti and F. Nitti
arXiv:0812.0792*

RG flow equations

$$Z = 1 - \frac{9}{4} X^2 + 2Y + 8Z + 8YZ + 4Z^2$$

$$\frac{dX}{d\phi} = -\frac{2}{9} Z \left(1 + \frac{1}{X} \frac{2\partial_\phi V - H_1^2 \partial_\phi f_1 + H_2^2 \partial_\phi f_2 - 2X H_2^2 f_2}{2V + H_1^2 f_1 + H_2^2 f_2} \right)$$

$$\frac{dY}{d\phi} = -\frac{2Y}{9X} Z \left(1 + \frac{1}{2Y} \frac{3H_1^2 f_1 - 4Y H_2^2 f_2}{2V + H_1^2 f_1 + H_2^2 f_2} \right)$$

$$\frac{dZ}{d\phi} = \frac{1-2Z}{9X} Z \left(1 + \frac{1+4Z}{1-2Z} \frac{H_2^2 f_2}{2V + H_1^2 f_1 + H_2^2 f_2} \right)$$

$$\frac{dH_1}{d\phi} = - \left(\frac{1+4Z}{3X} + \frac{\partial_\phi f_1}{f_1} \right) H_1$$

$$\frac{dH_2}{d\phi} = -\frac{4Z}{3X} H_2$$

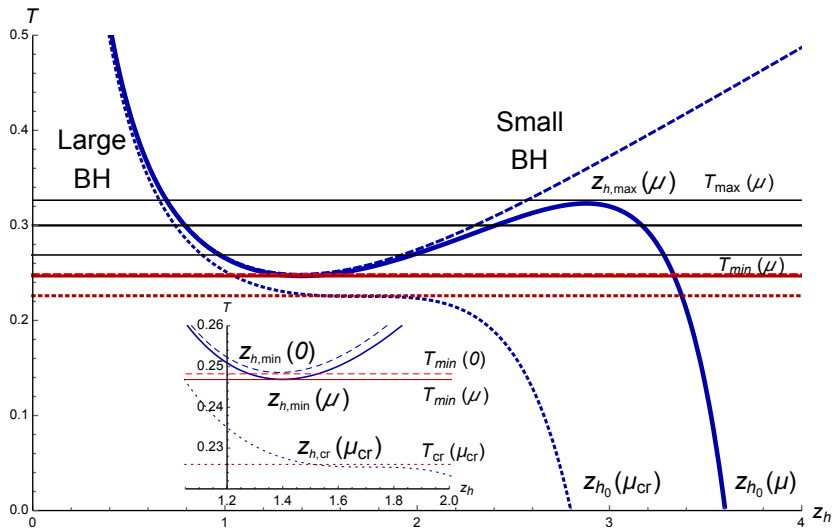
Thermodynamics for anisotropic solution

$$b(z) = \exp(cz^2/2), \nu \neq 1$$

$$s = \frac{e^{\frac{3cz_h^2}{4}}}{4z_h^{1+\frac{2}{\nu}}}$$

$$T = \frac{e^{-\frac{3cz_h^2}{4}}}{2\pi z_h} \left| \frac{1}{\mathfrak{G}\left(\frac{3}{4}cz_h^2\right)} + \frac{\mu^2 cz_h^{2+\frac{2}{\nu}} e^{\frac{cz_h^2}{4}}}{4 \left(1 - e^{\frac{cz_h^2}{4}}\right)^2} \left(1 - e^{\frac{cz_h^2}{4}} \frac{\mathfrak{G}(cz_h^2)}{\mathfrak{G}\left(\frac{3}{4}cz_h^2\right)}\right) \right|$$

Temperature $T(z_h, \mu)$ for $c = -1, \nu = 4.5$

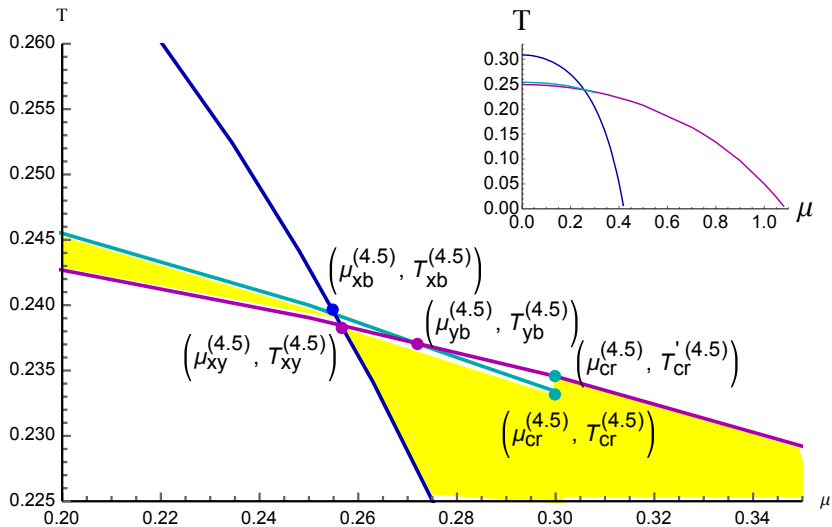


Confinement-deconfinement for $b(z) = \exp(cz^2/2)$, $\nu \neq 1$

$$DW_x \equiv cz + \frac{1}{\nu z} \sqrt{\frac{2}{3}} \sqrt{3c\nu^2 z^2 \left(\frac{cz^2}{2} - 3 \right) + 4\nu - 4} + \frac{g'}{2g} - \frac{2}{z} \Big|_{z=z_{DW_x}} = 0$$

$$DW_y \equiv cz + \frac{1}{\nu z} \sqrt{\frac{2}{3}} \sqrt{3c\nu^2 z^2 \left(\frac{cz^2}{2} - 3 \right) + 4\nu - 4} + \frac{g'}{2g} - \frac{\nu + 1}{\nu z} \Big|_{z=z_{DW_y}} = 0$$

Confinement-deconfinement phase transition for $c = -1, \nu = 4.5$



Conclusions

- The *AdS* anisotropic black hole solution for 5-dimensional Maxwell-dilaton gravity is found.
- The AGG solution is generalized for the non-zero chemical potential case.
- Scalar field restricts the BH radius and the warp-factor values.
- RG flow equations for anisotropic case are found.
- Confinement-deconfinement phase transition diagram is obtained [next talk]

To do:

- Investigate $P(z) \sim z^4$ to reconstruct Cornell potential.
- Shock-wave consideration
(Aref'eva, Mamedov, KR, Arifulov) – in progress.
- Drag-force consideration (Aref'eva, KR, Slepov) – in progress.

Thank you
for your attention